

1 stepped pressure equilibrium code : fu00aa

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1.1 outline

- Evaluates magnetic pressure and helicity volume integrals, and their derivatives with respect to interface geometry.

1.1.1 The magnetic pressure, $\int_V B^2 dv$, and the helicity, $\int_V \mathbf{A} \cdot \mathbf{B} dv$.

- In the annular regions, the volume integrals are given by the following:

$$\begin{aligned}
& \int_{s_{l-1}}^{s_l} ds \oint \oint d\theta d\zeta \sqrt{g} B^s B^s g_{ss} \\
= & \sum_{i=1}^{N_i} \sum_{j,k} \int_{s_{i-1}}^{s_i} ds (-m_j A_{\zeta,l,j} - n_j A_{\theta,l,j}) (-m_k A_{\zeta,l,k} - n_k A_{\theta,l,k}) \left\langle s_j \left| \sqrt{g}^{-1} g_{ss} \right| s_k \right\rangle \\
= & m_j A_{\zeta,l,j,p,i-1+l} m_k A_{\zeta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi_{r,q} \left\langle s_j \left| \sqrt{g}^{-1} g_{ss} \right| s_k \right\rangle \\
+ & m_j A_{\zeta,l,j,p,i-1+l} n_k A_{\theta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi_{r,q} \left\langle s_j \left| \sqrt{g}^{-1} g_{ss} \right| s_k \right\rangle \\
+ & n_j A_{\theta,l,j,p,i-1+l} m_k A_{\zeta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi_{r,q} \left\langle s_j \left| \sqrt{g}^{-1} g_{ss} \right| s_k \right\rangle \\
+ & n_j A_{\theta,l,j,p,i-1+l} n_k A_{\theta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi_{r,q} \left\langle s_j \left| \sqrt{g}^{-1} g_{ss} \right| s_k \right\rangle
\end{aligned} \tag{1}$$

$$\begin{aligned}
& \int_{s_{l-1}}^{s_l} ds \oint \oint d\theta d\zeta \sqrt{g} B^s B^\theta g_{s\theta} \\
= & \sum_{i=1}^{N_i} \sum_{j,k} \int_{s_{i-1}}^{s_i} ds (-m_j A_{\zeta,l,j} - n_j A_{\theta,l,j}) (-A'_{\zeta,l,k}) \left\langle s_j \left| \sqrt{g}^{-1} g_{s\theta} \right| c_k \right\rangle \\
= & m_j A_{\zeta,l,j,p,i-1+l} A_{\zeta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi'_{r,q} \left\langle s_j \left| \sqrt{g}^{-1} g_{s\theta} \right| c_k \right\rangle \\
+ & n_j A_{\theta,l,j,p,i-1+l} A_{\zeta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi'_{r,q} \left\langle s_j \left| \sqrt{g}^{-1} g_{s\theta} \right| c_k \right\rangle
\end{aligned} \tag{2}$$

$$\begin{aligned}
& \int_{s_{l-1}}^{s_l} ds \oint \oint d\theta d\zeta \sqrt{g} B^s B^\zeta g_{s\zeta} \\
= & \sum_{i=1}^{N_i} \sum_{j,k} \int_{s_{i-1}}^{s_i} ds (-m_j A_{\zeta,l,j} - n_j A_{\theta,l,j}) (A'_{\theta,l,k}) \left\langle s_j \left| \sqrt{g}^{-1} g_{s\zeta} \right| c_k \right\rangle \\
= & -m_j A_{\zeta,l,j,p,i-1+l} A_{\theta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi'_{r,q} \left\langle s_j \left| \sqrt{g}^{-1} g_{s\zeta} \right| c_k \right\rangle \\
+ & -n_j A_{\theta,l,j,p,i-1+l} A_{\theta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi'_{r,q} \left\langle s_j \left| \sqrt{g}^{-1} g_{s\zeta} \right| c_k \right\rangle
\end{aligned} \tag{3}$$

$$\begin{aligned}
& \int_{s_{l-1}}^{s_l} ds \oint \oint d\theta d\zeta \sqrt{g} B^\theta B^\theta g_{\theta\theta} \\
&= \sum_{i=1}^{N_i} \sum_{j,k} \int_{s_{i-1}}^{s_i} ds (-A'_{\zeta,l,j})(-A'_{\zeta,l,k}) \left\langle c_j \left| \sqrt{g^{-1}} g_{\theta\theta} \right| c_k \right\rangle \\
&= A_{\zeta,l,j,p,i-1+l} A_{\zeta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi'_{l,p} \varphi'_{r,q} \left\langle c_j \left| \sqrt{g^{-1}} g_{\theta\theta} \right| c_k \right\rangle
\end{aligned} \tag{4}$$

$$\begin{aligned}
& \int_{s_{l-1}}^{s_l} ds \oint \oint d\theta d\zeta \sqrt{g} B^\theta B^\zeta g_{\theta\zeta} \\
&= \sum_{i=1}^{N_i} \sum_{j,k} \int_{s_{i-1}}^{s_i} ds (-A'_{\zeta,l,j})(A'_{\theta,l,k}) \left\langle c_j \left| \sqrt{g^{-1}} g_{\theta\zeta} \right| c_k \right\rangle \\
&= -A_{\zeta,l,j,p,i-1+l} A_{\theta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi'_{l,p} \varphi'_{r,q} \left\langle c_j \left| \sqrt{g^{-1}} g_{\theta\zeta} \right| c_k \right\rangle
\end{aligned} \tag{5}$$

$$\begin{aligned}
& \int_{s_{l-1}}^{s_l} ds \oint \oint d\theta d\zeta \sqrt{g} B^\zeta B^\zeta g_{\zeta\zeta} \\
&= \sum_{i=1}^{N_i} \sum_{j,k} \int_{s_{i-1}}^{s_i} ds (A'_{\theta,l,j})(A'_{\theta,l,k}) \left\langle c_j \left| \sqrt{g^{-1}} g_{\zeta\zeta} \right| c_k \right\rangle \\
&= A_{\theta,l,j,p,i-1+l} A_{\theta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi'_{l,p} \varphi'_{r,q} \left\langle c_j \left| \sqrt{g^{-1}} g_{\zeta\zeta} \right| c_k \right\rangle
\end{aligned} \tag{6}$$

$$\begin{aligned}
& \int_{s_{l-1}}^{s_l} ds \oint \oint d\theta d\zeta \sqrt{g} \mathbf{A} \cdot \mathbf{B} \\
&= \sum_{i=1}^{N_i} \int_{s_{i-1}}^{s_i} ds \oint \oint d\theta d\zeta [-A_\theta \partial_s A_\zeta + A_\zeta \partial_s A_\theta] \\
&= \sum_{i=1}^{N_i} \sum_{j,k} \int_{s_{i-1}}^{s_i} ds \left[-A_{\theta,l,j} A'_{\zeta,l,k} + A_{\zeta,l,j} A'_{\theta,l,k} \right] \langle c_j | 1 | c_k \rangle \\
&= -A_{\theta,l,j,p,i-1+l} A_{\zeta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi'_{r,q} \langle c_j | 1 | c_k \rangle \\
&\quad + A_{\zeta,l,j,p,i-1+l} A_{\theta,l,k,q,i-1+r} \int_{s_{i-1}}^{s_i} ds \varphi_{l,p} \varphi'_{r,q} \langle c_j | 1 | c_k \rangle
\end{aligned} \tag{7}$$

3. In the innermost volume, there are additional radial factor $s^{m/2}$, but these are incorporated into the `igss`, `igst`, `igsz`, `igtt`, `igtz`, `igzz`, and `igoo`.

4. The above integrals are returned in

$$\text{1BBintegral(lvol)} = \int_V dv \mathbf{B} \cdot \mathbf{B}$$

$$\text{1ABintegral(lvol)} = \int_V dv \mathbf{A} \cdot \mathbf{B}$$

1.1.2 gauge invariant helicity

5. A gauge invariant form for the helicity is derived as follows. The change in helicity, δH , where the helicity is defined $H \equiv \int \mathbf{A} \cdot \mathbf{B} dv$, resulting from a change in the gauge, e.g. $\mathbf{A}' = \mathbf{A} + \nabla g$, is given by

$$\delta H = \int_V \nabla g \cdot \mathbf{B} dv = \int_V \nabla \cdot (g \mathbf{B}) dv = \int_{\partial V} g \mathbf{B} \cdot d\mathbf{s}. \tag{8}$$

6. The gauge function, g , is not required to be a single valued function of position, i.e. $g(s, \theta + 2\pi, \zeta) \neq g(s, \theta, \zeta)$ and $g(s, \theta, \zeta + 2\pi) \neq g(s, \theta, \zeta)$. In contrast, the magnetic field, and Jacobian, are periodic.
7. Using $\mathbf{B} \cdot \mathbf{n} = 0$ on the ideal interfaces that bound the volume, \mathcal{V} , the expression for the variation in helicity can be reduced to

$$\delta H = \iint [g(s, \theta, 2\pi) - g(s, \theta, 0)] \mathbf{B} \cdot d\mathbf{s} + \iint [g(s, 2\pi, \zeta) - g(s, 0, \zeta)] \mathbf{B} \cdot d\mathbf{s}, \quad (9)$$

(10)

where the first integral is over the region of the plane $\zeta = 0$ bounded by the adjacent interfaces, and the second integral over the region of the $\theta = 0$ plane.

8. HOW DO THE TERMS INVOLVING g GET PULLED OUT OF THE INTEGRALS?

Is $g(s, \theta, 2\pi) - g(s, \theta, 0) = \text{const.}$ and $g(s, 2\pi, \zeta) - g(s, 0, \zeta) = \text{const.}?$

9. If the vector potential is required to be single valued, the representation of the gauge function is restricted to be of the form

$$g = G_\theta \theta + G_\zeta \zeta + \tilde{g}(s, \theta, \zeta) \quad (11)$$

where G_θ and G_ζ are constants and \tilde{g} is periodic. In this case, $g(s, \theta, 2\pi) - g(s, \theta, 0) = 2\pi G_\zeta$ and $g(s, 2\pi, \zeta) - g(s, 0, \zeta) = 2\pi G_\theta$.

10. By identifying the toroidal and poloidal fluxes enclosed between the interfaces, $\delta\psi \equiv \iint \mathbf{B} \cdot d\mathbf{s}$, we obtain

$$\delta H = 2\pi G_\zeta \delta\psi_t + 2\pi G_\theta \delta\psi_p. \quad (12)$$

11. Some logic missing . . .

12. A gauge invariant form for the helicity is

$$H = \int \mathbf{A} \cdot \mathbf{B} dv - \delta\psi_p \int_0^{2\pi} A_\theta(s_{l-1}, \theta, \zeta) d\theta - \delta\psi_t \int_0^{2\pi} A_\zeta(s_l, \theta, \zeta) d\zeta, \quad (13)$$

where s_{l-1} coincides with the inner interface, and s_l coincides with the outer interface.

13. Note that the gauge and boundary conditions employed give $A_\theta(s_{l-1}, \theta, \zeta) = \psi_{t,l-1}$ and $A_\zeta(s_l, \theta, \zeta) = \psi_{p,l}$. After using $\delta\psi_t \equiv \psi_{t,l} - \psi_{t,l-1}$ and $\delta\psi_p \equiv \psi_{p,l} - \psi_{p,l-1}$, the final expression for the helicity is

$$H = \int \mathbf{A} \cdot \mathbf{B} dv + \psi_{t,l}\psi_{p,l} - \psi_{t,l-1}\psi_{p,l-1}. \quad (14)$$